Lieb-Thirring inequalities for eigenvalues of Schrödinger and Dirac operators

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27 Feb. 2016

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27 Feb. 2016, Tokyo

Schrödinger operators

Definition (Schrödinger operators)

Let

$$H = -\triangle + V(x)$$
 on $L^2(\mathbb{R}^d)$,

where
$$\triangle = \sum_{j=1}^{d} \frac{\partial^2}{\partial x_j^2}$$
 denotes the Laplace operator in \mathbb{R}^d
and V is a real-valued potential function of $x \in \mathbb{R}^d$.

Example (Hydrogenic atom)

$$H = - \bigtriangleup - rac{Z}{|x|}$$
 on $L^2(\mathbb{R}^d)$,

Our Goal

Energy eigenvalues for Schrödinger operators

 E_0, E_1, E_2, \ldots denote all non-positive eigenvalues of H.

The Riesz mean of order $\gamma \geq \mathbf{0}$

$$\sum_{j} |E_j|^{\gamma}$$

should be estimated in the view of V. In particular,

$$\sum_j |E_j|^\gamma = \left\{egin{array}{cc} ext{the number of eigenvalues} & ext{for } \gamma = 0 \ ext{the sum of possible energies} & ext{for } \gamma = 1. \end{array}
ight.$$

Lieb-Thirring inequalities

Theorem (Lieb and Thirring, 1976)

Let $\gamma \geq 0$. Assume that $V_{-}(x) = \max\{-V(x), 0\}$ satisfies the condition $V_{-} \in L^{\gamma+d/2}(\mathbb{R}^{d})$. Then, there is $L_{\gamma,d} > 0$ which is independent of V such that

$$\sum_{j} |E_{j}|^{\gamma} \leq L_{\gamma,d} \int_{\mathbb{R}^{d}} V_{-}(x)^{\gamma+d/2} dx \qquad (1)$$

holds when $\gamma \ge 1/2$ for d = 1, $\gamma > 0$ for d = 2, and $\gamma \ge 0$ for $d \ge 3$. Otherwise, there is V that violates the inequalities (4) for any finite choice of $L_{\gamma,d}$.

Remarks on Lieb-Thirring inequalities

- Cwikel [1], Lieb [7] and Rozenbljum [10]; the critical case $\gamma = 0$ for $d \ge 3$. (1970s)
- Lieb and Thirring [9]; almost all cases in Lieb-Thirring inequalities. (1976)
- Weidl [13]; the remaining case $\gamma = 1/2$ for d = 1. (1996)

Semi-classical approximation

The coefficients $L_{\gamma,d}$ should be compared to classical ones obtained by the semi-classical approximation

$$L^{\mathsf{cl}}_{\gamma,d} = (2\pi)^d \int_{|\pmb{p}| \leq 1} (1-|\pmb{p}|^2)^\gamma \; d\pmb{p} = rac{\Gamma(\gamma+1)}{(4\pi)^{d/2} \Gamma(\gamma+1+d/2)}$$

where Γ is the Gamma function. It is also known that $L_{\gamma,d}/L_{\gamma,d}^{cl} \ge 1$ for all possible γ and d, and that $L_{\gamma,d}/L_{\gamma,d}^{cl}$ is non-increasing on γ . (Aizenman and Lieb, 1978)

Semi-classical approximation

- Helffer and Robert [3]; $L_{\gamma,d}/L_{\gamma,d}^{\mathsf{cl}}>1$ for $\gamma<1.$ (1990)
- Hundertmark, Lieb and Thomas [5]; L_{1/2,1} = 2L^{cl}_{1/2,1}. (1998)
- Hundertmark, Laptev and Weidl [4]; $L_{\gamma,1} = L_{\gamma,1}^{cl}$ for $\gamma \ge 3/2$. (2000)
- Laptev and Weidl [6] enables that L_{γ,d} = L^{cl}_{γ,d} if γ ≥ 3/2 for all d. (2000)

Recent results (1)

Dolbeault, Laptev and Loss [2] have improved the coefficient which is also known as best possible at the present time.

Dolbeault, Laptev and Loss, 2008

$$L_{1,d}/L_{1,d}^{
m cl} \le \pi/\sqrt{3} = 1.81...$$
 for all d .

[2] J. Dolbeault, A. Laptev and M. Loss; J. Eur. Math. Soc., Vol. 10, pp. 1121-1126. (2008)

Recent results (2)

Rumin $\left[11\right]$ and Solovej $\left[12\right]$ has proposed a new approach of proving that

Rumin and Solovej, 2011

$$L_{1,d}/L_{1,d}^{\mathsf{cl}} \leq \left(rac{d+4}{d}
ight)^{d/2}$$

[11] M. Rumin, Duke Math. J., 160, no. 3, 567–597. (2011)
[12] J. P. Solovej,
"The Lieb-Thirring inequality." (2011)

Lieb-Thirring conjecture

It is conjectured by Lieb and Thirring [9] that the optimal $L_{1,3}$ coincides with $L_{1,3}^{cl}$.

Conjecture (Lieb and Thirring, 1976)

$$L_{1,3} = L_{1,3}^{cl}$$

[9] E. H. Lieb and W. E. Thirring, "Inequalities for the Moments of the Eigenvalues of the Schrödinger Hamiltonian and their Relation to Sobolev Inequarities," in *Studies in Mathematical Physics* (1976), pp. 269–303.

Stability of Matter

$$H = \frac{1}{2} \sum_{j=1}^{N} \left(-i \nabla_j + \sqrt{\alpha} A(x_j) \right)^2 + \alpha V(X, R),$$

where $\alpha = e^2/\hbar c \approx 1/137 > 0$ is Sommerfeld's fine structure constant, A is an arbitrary magnetic vector potential in $L^2_{loc}(\mathbb{R}^3; \mathbb{R}^3)$. The Coulomb potential is written by

$$V(X,R) = \sum_{j=1}^{N} \sum_{k=j+1}^{N} \frac{1}{|x_j - x_k|} - \sum_{j=1}^{N} \sum_{k=1}^{M} \frac{Z_k}{|x_j - R_k|} + \sum_{j=1}^{M} \sum_{k=j+1}^{M} \frac{Z_j Z_k}{|R_j - R_k|},$$

where M is number of nuclei.

Stability of Matter Why is our world stable?

Lieb and Thirring [9] have improved the result by Dyson and Lenard for the stability of non-relativistic matter.

Theorem (Stability of matter of the second kind)

Let $Z_{max} = \max_{j} \{Z_j\}$. For all normalized, antisymmetric wavefunction ψ with q spin states,

$$(\psi, H\psi) \ge -0.231 \alpha^2 N q^{2/3} \left(1 + 2.16 Z_{max} (M/N)^{1/3}\right)^2$$

It should be remarked that $N^{1/3}M^{2/3} \le N + M$ implies the linear dependence on the total number of particles.

Pseudo-relativistic operators

Definition (Pseudo-relativistic operators)

Let

$$H = |-i\nabla| + V(x)$$
 on $L^2(\mathbb{R}^d)$,

where ∇ denotes the differential vector in \mathbb{R}^d and V is a real-valued potential function of $x \in \mathbb{R}^d$.

It should be noted for relativistic mechanics that it is necessary to consider the operators

$$T = \sqrt{\mathbf{p}^2 + \mu^2} - \mu \tag{2}$$

for $\mu > 0$, where

$$\sqrt{-\Delta} - \mu \le \sqrt{\mathbf{p}^2 + \mu^2} - \mu \le \sqrt{-\Delta}.$$
 (3)

LT inequalities for pseudo-relativistic operators

Theorem (Lieb 1980, and Daubechies 1983)

Let $\gamma \geq 0$. Assume that $V_{-}(x) = \max\{-V(x), 0\}$ satisfies the condition $V_{-} \in L^{\gamma+d}(\mathbb{R}^{d})$. Then, there is $L_{\gamma,d} > 0$ which is independent of V such that

$$\sum_{j} |E_{j}|^{\gamma} \leq L_{\gamma,d} \int_{\mathbb{R}^{d}} V_{-}(x)^{\gamma+d} dx$$
(4)

holds when $\gamma > 0$ for d = 1, and $\gamma \ge 0$ for $d \ge 2$.

Dirac operators

Definition (Dirac operators, d = 3)

Let

$$H_D = \alpha \cdot \mathbf{p} + \beta + V(x)$$
 on $L^2(\mathbb{R}^3)$,

where $\mathbf{p} = -i\nabla$ denotes the momentum operator in \mathbb{R}^3 and V is a real-valued potential function of $x \in \mathbb{R}^3$. Here,

$$\alpha = \begin{pmatrix} 0 & \sigma \\ \sigma & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} l_2 & 0 \\ 0 & -l_2 \end{pmatrix}$$
(5)

where $\sigma = (\sigma_1, \sigma_2, \sigma_3)$ are the three Pauli matrices in standard representation:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$
(6)

Lieb-Thirring inequalities for Dirac operators (1)

Theorem (Cancelier, Lévy-Bruhl and Nourrhgat, 1995)

Let $\gamma \geq 0.$ Then, there is $C_{\gamma} > 0$ which is independent of V such that

$$\sum_{j} (1 - E_j^2)^{\gamma} \leq C_{\gamma} \int_{\mathbb{R}^d} |V(x)|^{\gamma} (|V(x)|^3 + |V(x)|^{3/2}) \, dx.$$
 (7)

The proof of the Theorem follows from

$$\sum_{j} (1 - E_j^2)^{\gamma} = \gamma \int_0^1 \lambda^{\gamma - 1} \mathcal{N}(E, V) \ dE, \qquad (8)$$

where N(E, V) is the number of eigenvalues of H_D in $(-\sqrt{1-E}, \sqrt{1-E})$.

The number of eigenvalues for Dirac operators (2)

$$V \in L^{p}(\mathbb{R}^{3}) + \varepsilon L^{\infty}(\mathbb{R}^{3}) \Leftrightarrow \int_{|V(x)| \ge \varepsilon} |V(x)|^{p} dx < \infty$$
 (9)

Theorem (Cancelier, Lévy-Bruhl and Nourrhgat, 1995)

Assume that V(x) satisfies the condition $V \in L^{3}(\mathbb{R}^{3}) + \frac{E}{4}L^{\infty}(\mathbb{R}^{3})$. Then, there is C > 0 which is independent of V such that

$$N(E,V) \leq C \int_{|V(x)| \geq E/4} (|V(x)|^3 + |V(x)|^{3/2}) dx.$$
 (10)

The finite number of eigenvalues

Corollary (Cancelier, Lévy-Bruhl and Nourrhgat, 1995)

If $V \in L^3(\mathbb{R}^3) \cap L^{3/2}(\mathbb{R}^3)$, the number N(V) of eigenvalues of H_D in (-1, 1) is finite, and there is C > 0 which is independent of V such that

$$N(V) \leq C \int_{\mathbb{R}^d} (|V(x)|^3 + |V(x)|^{3/2}) dx.$$
 (11)

Concluding remarks

- We reviewed the topics around Lieb-Thirring inequalities for Schrödinger operators, pseudo-relativistic operators and Dirac operators.
- We share the proof of Lieb-Thirring inequalities for Dirac operators with rough value of coefficients..
- Lieb-Thirring conjecture remains open at the present time.